

## Singlet Oxygen Chemistry in Water. 2. Photoexcited Sensitizer Quenching by O<sub>2</sub> at the Water–Porous Glass Interface

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Insight into the O<sub>2</sub> quenching mechanism of a photosensitizer (static or dynamic) would be useful for the design of heterogeneous systems to control the mode of generation of <sup>1</sup>O<sub>2</sub> in water. Here, we describe the use of a photosensitizer, *meso*-tetra(*N*-methyl-4-pyridyl)porphine (**1**), which was adsorbed onto porous Vycor glass (PVG). A maximum loading of  $1.1 \times 10^{-6}$  mol **1** per g PVG was achieved. Less than 1% of the PVG surface was covered with photosensitizer **1**, and the penetration of **1** reaches a depth of 0.32 mm along all faces of the glass. Time-resolved measurements showed that the lifetime of triplet **1**\*-ads was 57 μs in water. Triplet O<sub>2</sub> quenched the transient absorption of triplet **1**\*-ads; for samples containing  $0.9 \times 10^{-6}$ – $0.9 \times 10^{-8}$  mol **1** adsorbed per g PVG, the Stern–Volmer constant, *K*<sub>D</sub>, ranged from 23 700 to 32 100 M<sup>-1</sup>. The adduct formation constant, *K*<sub>S</sub>, ranged from 1310 to 510 M<sup>-1</sup>. The amplitude of the absorption at 470 nm decreased slightly (by about 0.1) with increased O<sub>2</sub> concentrations. Thus, the quenching behavior of triplet **1**\*-ads by O<sub>2</sub> was proposed to be strongly dependent on dynamic quenching. Only ~10% of the quenching was attributed to the static quenching mechanism. The quenching of triplet **1**\*-ads was similar to that observed for photosensitizers in homogeneous solution which are often quenched dynamically by O<sub>2</sub>.

### 1. Introduction

Heterogeneous materials have been studied for many years, and chemists discovered that some could be used as supports for the generation of singlet oxygen [<sup>1</sup>O<sub>2</sub>(<sup>1</sup>Δ<sub>g</sub>)].<sup>1–1c</sup> Recently, we reported that *meso*-tetra(*N*-methyl-4-pyridyl)porphine (**1**) adsorbed onto porous Vycor glass (PVG), in which a pH decrease of the surrounding solution indicated the displacement of protons from the surface silanol groups via cation exchange (Scheme 1).<sup>2</sup> Singlet oxygen was generated cleanly in aqueous solution upon irradiation of this heterogeneous sensitizer, **1**-ads.<sup>2</sup>

Despite the effectiveness of this and other heterogeneous systems to generate <sup>1</sup>O<sub>2</sub>, surprisingly little is known about the mechanism of sensitizer quenching by O<sub>2</sub> at *water–solid* interfaces. For example, how does the oxygen encounter the excited PVG heterogeneous sensitizer? What mechanism (static or dynamic) converts ground-state O<sub>2</sub> into <sup>1</sup>O<sub>2</sub>, which then diffuses into the bulk solution?

Some detail of the O<sub>2</sub> quenching process may be gleaned from previous studies of *gas–solid* systems. These studies often indicate the static quenching of O<sub>2</sub>—where a ground-state O<sub>2</sub> adduct is formed at the surface—rather than a dynamic encounter of O<sub>2</sub> with the surface. At the *gas–solid* interface, Gafney showed that photoexcited Ru(bpy)<sub>3</sub><sup>2+</sup>-adsorbed to PVG statically quenches O<sub>2</sub>.<sup>3</sup> Avnir showed that photoexcited Ru(bpy)<sub>3</sub><sup>2+</sup>-adsorbed to porous silica and porous glass statically quench O<sub>2</sub> at low temperatures, but dynamically quench O<sub>2</sub> at high temperatures,<sup>4</sup> and Thomas showed that aromatic compounds adsorbed onto nonporous glass statically quench O<sub>2</sub>.<sup>5,6</sup> Pore size

of silica-adsorbed photosensitizers can be tailored to statically or dynamically quench O<sub>2</sub> in *gas–solid* systems.<sup>7,8</sup>

Few studies have examined sensitizer quenching by O<sub>2</sub> at organic solvent–solid interfaces<sup>9</sup> or water–solid interfaces.<sup>10</sup> Despite reports on <sup>1</sup>O<sub>2</sub> production (Φ<sub>Δ</sub>) and <sup>1</sup>O<sub>2</sub> lifetimes (τ<sub>Δ</sub>) in homogeneous aqueous media, Nafion membranes,<sup>11</sup> micellar media,<sup>12</sup> water-soluble supramolecular hosts,<sup>13</sup> or aqueous reactions using TiO<sub>2</sub>,<sup>14</sup> studies of the quenching mechanism of heterogeneous photosensitizers by O<sub>2</sub> at the *water–solid* interface are uncommon and are in need of elaboration.

We report here a photophysical analysis of how O<sub>2</sub> quenches **31**\* at the *water–PVG* interface. We also examined the effects of surface loading and percent coverage of the photosensitizer and the depth that **1** can penetrate into the PVG material. An assessment of the quenching of the photosensitizer (static or dynamic) could help in the design of heterogeneous systems that attempt to control the exact mode of generation of <sup>1</sup>O<sub>2</sub> in water.

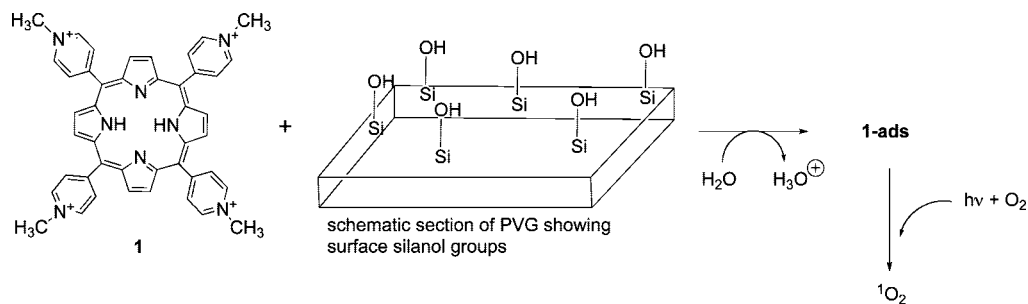
### 2. Experimental Section

**2.1. Materials and Sample Preparation.** Deionized H<sub>2</sub>O was obtained from a U.S. Filter Corp. deionization system. *meso*-Tetra(*N*-methyl-4-pyridyl)porphine tetratosylate was purchased commercially and used as received. The PVG was used as a host material and was purchased from Advanced Glass and Ceramics (Corning 7930). The PVG has a void space of ~28% of the volume, average pore sizes of 40 Å, a hydrophilic absorbing surface area of 250 m<sup>2</sup>/g, a density of 1.38 g/mL, and transparency in the near-UV (50% T at 351 nm), visible, and parts of the near-IR.<sup>15</sup> Pieces of PVG (1.5 cm × 1.5 cm; thickness: 1.15–1.53 mm) were heated in a muffle furnace at 500 °C and stored in a desiccator under vacuum at 30 mmHg.

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SCHEME 1: Adsorption of Photosensitizer **1** onto the PVG Surface and Production of Singlet Oxygen

The photosensitizer-coated PVG was prepared by soaking 1.4 g of PVG into a 20 mL  $3.7 \times 10^{-5}$  M aqueous solution of **1** for 15–64 h. The amount of **1** adsorbed onto PVG was calculated from the difference in absorbance of the solution before introduction of PVG and the absorbance of the same solution after the removal of PVG.<sup>2,16</sup> For the transient absorption studies, PVG was cut such that each piece fits diagonally into a 1 cm path-length quartz cuvette. Each PVG/**1** sample was placed into the quartz cuvette, which contained deionized water, and was sparged with  $\text{N}_2$  gas for 20 min. Any bubbles that stuck onto the PVG surface were removed by sonication for a few seconds. Some PVG samples (1.5 cm  $\times$  1.5 cm; thickness: 1.53 mm) were cut into 0.3 cm  $\times$  1.5 cm pieces and photographed with a microscope.

**2.2. Instruments.** Photographic images were taken with a Nikon TE200 microscope equipped with an Orca 100 monochrome charge-coupled device (CCD) camera and a Hamamatsu camera controller (C4742-95). The light source used was a 100 W mercury arc lamp. The objective used was a Plan Apo 60 $\times$ Oil DIC. Images were recorded and analyzed with CamMedia software.<sup>17</sup> Absorption spectra were collected with a Hitachi UV-vis U-2001, a Hewlett-Packard 8453 diode array, or a Varian Carey 14 spectrophotometer. In some experiments with the Carey 14 instrument, the Q-bands of **1** at 525, 552, 585, and 640 nm were followed because the intense Soret band at 422 nm gave absorbances of 2.5 or greater, which saturated the UV-vis detector. A Quantel Brilliant B Nd:YAG laser (417 nm, 0.6–1.6 mJ/pulse) was used in the transient absorption experiments.<sup>18,19</sup>

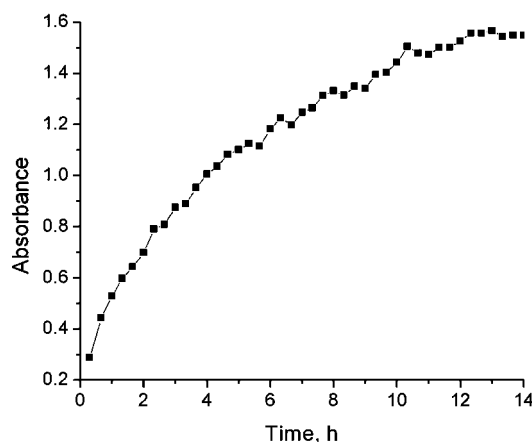
**2.3. Transient Absorption Spectroscopy.** The PVG samples used in these experiments contained  $0.9 \times 10^{-8}$ – $1.1 \times 10^{-6}$  mol **1**/g PVG. Room-temperature time-resolved measurements were conducted as previously described.<sup>20</sup> The PVG samples were excited at 417 nm but oriented so that reflected laser light was directed at a 45° angle away from the detector on the transient absorption apparatus. Each kinetic trace is an average of 64 laser pulses. The triplet of **1-ads** was monitored at 470 nm, in which a 435 nm long pass filter was placed in front of the detector monochromator to block scattered laser light. The data points in the transient absorption experiments were collected every 10 nm from 450 to 550 nm. An MKS Instruments multigas controller (MKS 647C) was used to control the flow of  $\text{O}_2$  and  $\text{N}_2$  through two MKS Instruments flow controllers. The flow controllers were set for a flow rate of 100 standard cubic centimeters per minute (sccm). Gas correction factors (GCFs) programmed into the MGC were utilized to correct for specific heat, density, and molecular structure of  $\text{O}_2$  and  $\text{N}_2$ . The samples were purged with different ratios of an  $\text{O}_2$ : $\text{N}_2$  stream of gas for 5 min before each kinetic trace was taken. The oxygen concentration of a solution bubbled with the  $\text{O}_2$ : $\text{N}_2$  gas stream were determined by using Henry's law.<sup>21,22</sup>

**2.4. Computational Methods.** Density function theoretical (DFT) calculations were conducted by the exchange-correlation of B3LYP along with Pople basis set 6-31G(d) with the use of the Gaussian 03 program package.<sup>23</sup> The solvent accessible surface was computed by the method of Lee and Richards.<sup>24</sup>

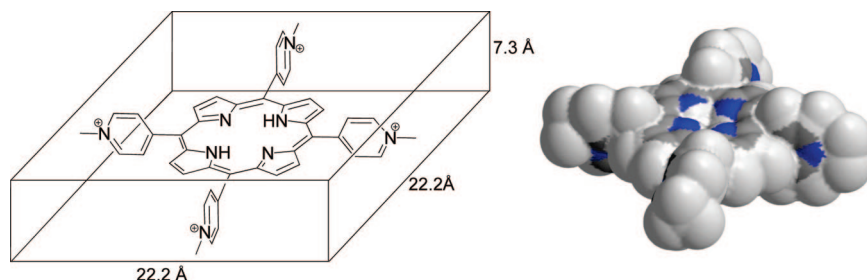
## 3. Results and Discussion

**3.1. Time-Dependent Adsorption of Photosensitizer onto PVG.** A clean piece of PVG placed into an aqueous solution containing **1** led to the adsorption of **1** onto PVG. Figure 1 shows the time-dependent adsorption of **1** onto PVG over a 15 h period. The adsorption process was followed by monitoring the largest of the four Q-bands of **1-ads** at  $\lambda = 525$  nm. A plateau is reached when there is  $8.8 \times 10^{-7}$  mol **1** adsorbed onto 1 g PVG. Loadings of  $1.1 \times 10^{-6}$  mol **1** onto PVG can be achieved, but only after a 48–72 h period. Once the **1-ads** samples contained the desired surface coverage, they were rinsed with distilled water prior to use.

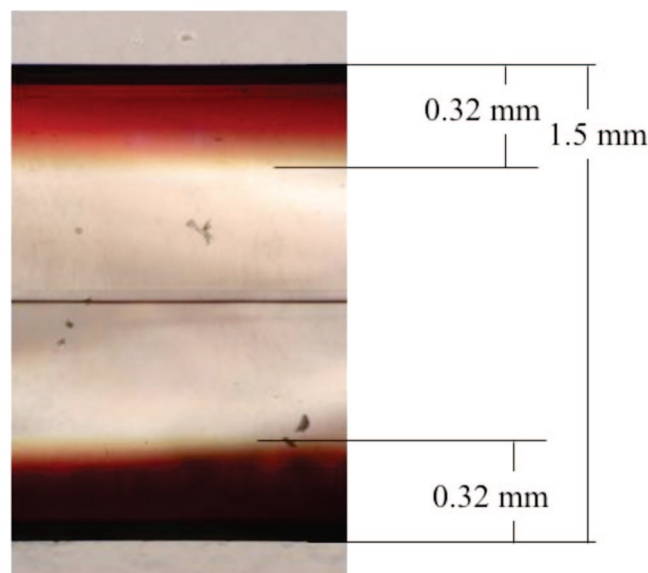
**3.2. Uniform Photosensitizer Distribution.** Absorption spectra were recorded at different points on the PVG film to analyze the dispersal of adsorbed **1**. The Q-band absorptions (475–700 nm) are found to be identical to within 0.01 absorbance unit, suggesting that the distribution of **1** is uniform; thus, the surface coverage of **1** adsorbed onto PVG could then be estimated. Two methods were used to determine the surface coverage of **1** adsorbed onto PVG. Both are in qualitative agreement with each other: (1) The amount of uncoated PVG was determined by subtracting 1.0 g of PVG from ( $8.83 \times 10^{-7}$  mol **1**/g PVG  $\times$  679.61 g/mol **1**), which is 0.9994 g. Dividing  $8.83 \times 10^{-7}$  mol **1**/g PVG into (0.9994 g uncoated PVG  $\times$  250 m<sup>2</sup>/g PVG) yielded  $3.5 \times 10^{-9}$  mol **1**/m<sup>2</sup>, indicating that



**Figure 1.** Adsorption of a PVG sample at  $\lambda = 525$  nm that was sitting in a solution containing **1**. The PVG sample was removed at the indicated times. The plateau region at 14 h corresponds to  $8.83 \times 10^{-7}$  mol **1** adsorbed onto 1 g PVG.



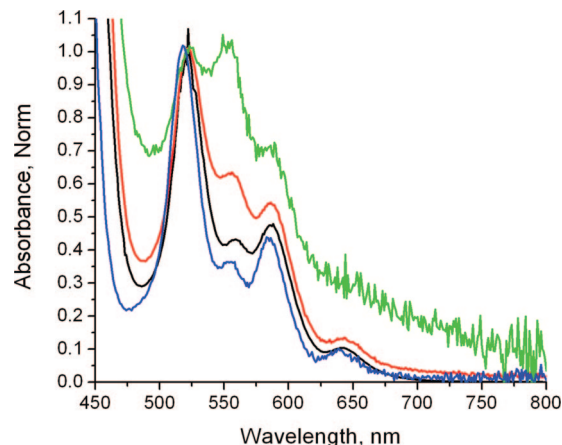
**Figure 2.** Rectangular shape of **1** (left) estimated from a B3LYP/6-31G(d) optimized structure and the corresponding computed solvent accessible map (right).



**Figure 3.** A low-magnification (60 $\times$ ) cross-sectional optical image. The red areas indicate the depth of **1** accessed into PVG. The image shows 0.32 mm penetration depth on each face for  $0.9 \times 10^{-6}$  mol **1** adsorbed onto a 1.5 mm PVG sample. The horizontal line in the middle of the PVG sample corresponds to where two glass layers meet due to the commercial fabrication process.

0.35% of the PVG surface is covered by **1**. (2) Surface coverage was determined by a calculation, in which porphyrin **1** is taken as a rectangular shape (22.2 Å  $\times$  22.2 Å  $\times$  7.3 Å) multiplied by ( $8.83 \times 10^{-7}$  mol **1**/g PVG  $\times$   $6.02 \times 10^{23}$  molecules/mol  $\times$   $10^{-21}$  mm<sup>3</sup>/Å<sup>3</sup>), which equals 1.91 mm<sup>3</sup>/g PVG. The rectangular shape of **1** was estimated on the basis of the B3LYP/6-31G(d) optimized structure and the corresponding solvent accessible contour map (Figure 2), in which the pyridinium rings are nearly orthogonal to the plane of the porphyrin (67°). The corresponding weight of **1**-adsorbed PVG (1.91 mm<sup>3</sup>  $\times$  1.38 g PVG/mL) equals 2.64 g PVG. Therefore, the surface coverage of **1** onto PVG was calculated by taking [ $8.83 \times 10^{-7}$  mol **1**/(2.64 g PVG  $\times$  250 m<sup>2</sup>/g PVG)] or  $1.4 \times 10^{-9}$  mol **1**/m<sup>2</sup>, yielding 0.13% PVG surface coverage by **1**. Both calculations [(1) and (2), vide supra] indicated that <1% of the PVG surface is covered with the photosensitizer **1**. For comparison, the surface coverage of Ru(bpy)<sub>3</sub><sup>2+</sup> on PVG is <1%,<sup>3</sup> on porous silica is ~6%, and on porous glass is ~10%.<sup>4</sup>

**3.3. Photosensitizer Penetration Depth.** The depth that **1** can penetrate into PVG was examined using a microscope equipped with a CCD camera. Figure 3 shows a 1.5 mm thick (sensitizer coated) PVG sample, cut so that the depth of **1** penetrated into PVG could be viewed. The microscope image shows the penetration of **1** reaches a maximum depth of 0.32 mm along all faces of the sample, which corresponds to the plateau region of  $8.8 \times 10^{-7}$  mol **1**/g PVG (Figure 1). Although

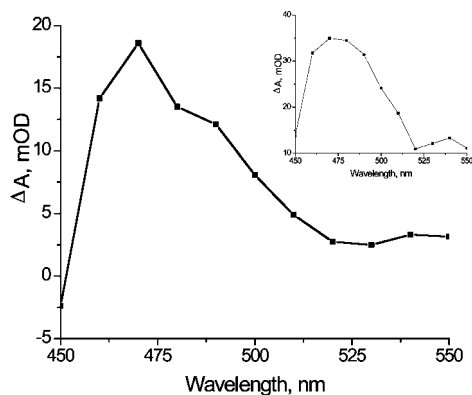


**Figure 4.** Absorption spectra of **1** in water (blue line) and **1-ads** at different loadings:  $0.9 \times 10^{-8}$  (green line),  $0.9 \times 10^{-7}$  (red line), and  $0.9 \times 10^{-6}$  (black line) mol of **1** adsorbed on PVG. The spectra were normalized at 532 nm. Except for the green line ( $0.9 \times 10^{-8}$  mol **1** adsorbed onto PVG), the Q-band maximum is at 525 nm.

O<sub>2</sub> and other gases are permeable and can pass through the connected pores of PVG,<sup>25</sup> **1** neither penetrates to the center of the PVG film nor is **1** localized on the outer surface of PVG. A 10-fold mole increase of **1** resulted in only a 4-fold local increase in sensitizer distribution into PVG [cf.  $0.9 \times 10^{-6}$  mol **1** (penetration depth 0.32 mm) and  $0.9 \times 10^{-7}$  mol **1** adsorbed onto 1 g PVG (penetration depth 0.09 mm)]. Oxygen is expected to reach the excited sites of **1-ads**, controlled by Knudsen diffusion, in which O<sub>2</sub> collides numerous times within the pore walls, eventually proceeding through the PVG channels. For comparison, Ru(bpy)<sub>3</sub><sup>2+</sup> penetrates  $0.5 \pm 0.1$  mm into PVG.<sup>3</sup> Streptocyanine dyes also possessed diffusivity into silica gels, influenced by the gel porosities.<sup>26</sup>

**3.4. Spectral Properties.** Previously, we reported that the spectral features of **1-ads** are nearly identical to **1** in fluid water solution.<sup>2</sup> Here, we report the spectral features measured at different concentrations of **1** adsorbed onto PVG.

Figure 4 shows absorption spectra of **1** in water and **1** adsorbed onto PVG. While the effect of higher coverage dose of **1** produces neither red shifts nor blue shifts in the Q-bands, the adsorption of **1** on PVG leads to an unexpected effect on the absorption intensity. Higher coverages of **1** adsorbed onto PVG led to a decrease of the Q-band absorption intensity, (cf.  $0.9 \times 10^{-8}$ ,  $0.9 \times 10^{-7}$ , and  $0.9 \times 10^{-6}$  mol **1**/g PVG). At higher surface coverage, the absorption spectrum of **1-ads** becomes somewhat similar to that of **1** in aqueous solution (cf. spectra in blue and black, Figure 4). Deducing a possible orientational effect of the porphyrin on the PVG surface is challenging based upon these normalized absorption spectra. Polarization effects are known for porphyrins result in symmetric changes of all the S<sub>0</sub>  $\rightarrow$  S<sub>1</sub> transitions. However, atomic-level



**Figure 5.** Nanosecond transient absorption spectra observed at 500 ns after pulsed-light excitation (417 nm, 1.6 mJ/pulse) of **1** in water solution (inset) and **1** at the water/PVG interface. Data points were collected every 10 nm from 450 to 550 nm.

detail is unavailable for PVG and how it may relate to the conformation of adsorbed **1**. Interestingly, recent studies have shown a reduced hydrating ability of  $\sim 2.2$  H-bonds per water molecule confined in PVG pores compared bulk water with  $\sim 3.6$  H-bonds per water molecule, pointing to the importance for confinement effects.<sup>27,28</sup>

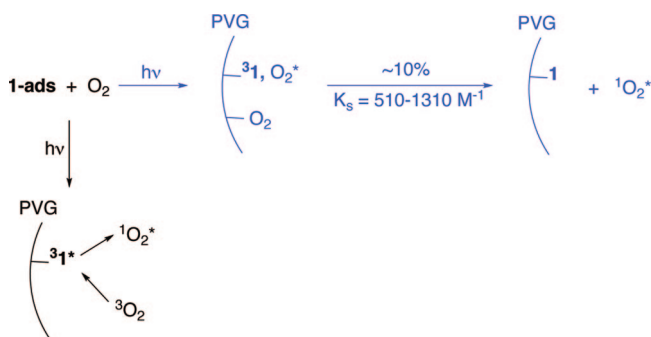
**3.5. Time-Resolved Photophysical Studies.** Figure 5 shows a nanosecond transient absorption spectrum generated from pulsed 417 nm light excitation of  $0.9 \times 10^{-7}$  mol **1** adsorbed onto 1 g PVG in an  $N_2$ -purged  $H_2O$  solution. The absorption band at  $\sim 470$  nm was assigned to  $^3\mathbf{1}^*\text{-ads}$ . Similar transient features were observed for  $^3\mathbf{1}^*$  in fluid water (Figure 5, inset) and also by Reddi et al. for  $^3\mathbf{1}^*$  in phosphate buffered solution and in aqueous 2% sodium dodecyl sulfate solution.<sup>29</sup> First-order decay kinetics were observed for the transient absorption of  $^3\mathbf{1}^*\text{-ads}$ , which decayed cleanly to baseline. The lifetime of  $^3\mathbf{1}^*\text{-ads}$  ( $\tau_0 = 57 \pm 1 \mu\text{s}$ ) is similar to  $^3\mathbf{1}^*$  in fluid aqueous solution ( $\tau_0 = 49 \pm 1 \mu\text{s}$ ). Our absorption measurements in the UV–vis region before and after transient absorption experiments revealed no significant change. Next we describe how  $O_2$  quenches the transient absorption of  $^3\mathbf{1}^*\text{-ads}$  and  $^3\mathbf{1}^*$  in water solution.

Figure S1 shows the lifetime and the amplitude quenching of the  $^3\mathbf{1}^*\text{-ads}$  transient absorption by  $O_2$  (Supporting Information). Photoexcited **1** quenching by  $O_2$  at the PVG/water interface encouraged a Stern–Volmer analysis for insight into quenching mechanism. Two mechanisms were considered: (1) the dynamic encounter of  $O_2$  with an excited site on the surface, and (2) a ground-state  $O_2$  adduct formed at the surface with migration of adsorbed  $O_2$  to an excited site (Scheme 2).

Stern–Volmer data collected at nine  $O_2$  concentrations (from 0 to 0.4 mM) are shown in Table 1 and Figure 6, in which the symbol  $\tau$  refers to the lifetime of the excited porphyrin,  $k_q$  is the rate constant for  $O_2$  quenching,  $K_D$  is the Stern–Volmer constant,  $K_S$  is the association constant for complex formation between the sensitizer and  $O_2$ , and the subscript “0” indicates data in the absence of  $O_2$ . The left axis of Figure 6 shows the plot of  $\tau_0/\tau$  vs  $[O_2]$  to be linear over the  $O_2$  concentration examined and corresponded to the Stern–Volmer equation:  $\tau_0/\tau = 1 + k_q\tau_0[O_2] = 1 + K_D[O_2]$ . For  $0.9 \times 10^{-7}$  mol **1** adsorbed onto 1 g PVG,  $K_D = 26\,500 \text{ M}^{-1}$  and the quenching rate constant  $k_q = 4.6 \times 10^8 \text{ M}^{-1} \text{ s}^{-1}$ , indicating that 50% of  $^3\mathbf{1}^*\text{-ads}$  was quenched at an  $[O_2]$  of 0.04 mM. For the samples containing  $0.9 \times 10^{-6}$ – $0.9 \times 10^{-8}$  mol **1** adsorbed onto 1 g PVG, the  $K_D$  values ranged from 32 000 to 23 700  $\text{M}^{-1}$ .

Interestingly, the bimolecular quenching constant  $k_q$  for **1-ads** ( $k_q = \sim 5 \times 10^8 \text{ M}^{-1} \text{ s}^{-1}$ ) is about one-quarter the

## SCHEME 2: Mechanism of $^3\mathbf{1}^*\text{-ads}$ Quenching by $O_2^a$

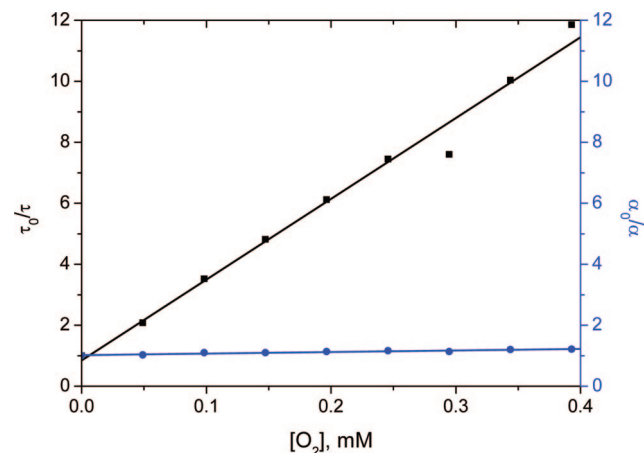


<sup>a</sup> Dynamic quenching route in black and static quenching route in blue.

**TABLE 1: Stern–Volmer Quenching Constants as a Function of Coverage of **1** onto Porous Vycor Glass or of **1** in Fluid Water Solution**

photosens	absorbance $\lambda_{520}$	$K_D$ ( $\text{M}^{-1}$ )	$K_S$ ( $\text{M}^{-1}$ )	$k_q$ ( $\text{M}^{-1} \text{ s}^{-1}$ )
<b>1-ads</b> <sup>a,b</sup>	0.87	$23700 \pm 1320$	$1310 \pm 210$	$4.2 \times 10^8$
<b>1-ads</b> <sup>a,c</sup>	0.25	$26500 \pm 1220$	$510 \pm 60$	$4.6 \times 10^8$
<b>1-ads</b> <sup>a,d</sup>	0.025	$32100 \pm 1460$	$650 \pm 80$	$5.6 \times 10^8$
<b>1</b> in fluid $H_2O$ <sup>e</sup>	$\sim 0.2$	$89600 \pm 2600$	$1160 \pm 170$	$1.8 \times 10^9$

<sup>a</sup> The PVG/**1** samples were immersed in water. <sup>b</sup> Sample contained  $0.9 \times 10^{-6}$  mol **1** adsorbed onto PVG. <sup>c</sup> Sample contained  $0.9 \times 10^{-7}$  mol **1** adsorbed onto PVG. <sup>d</sup> Sample contained  $0.9 \times 10^{-8}$  mol **1** adsorbed onto PVG. <sup>e</sup> Homogeneous **1** in deionized water solution.



**Figure 6.** Lifetime quenching as a function of  $O_2$  concentration of a  $H_2O$  solution containing  $0.9 \times 10^{-7}$  mol **1** adsorbed onto PVG (in black, left). Amplitude quenching as a function of  $O_2$  concentration of a  $H_2O$  solution containing  $0.9 \times 10^{-7}$  mol **1** adsorbed onto PVG (in blue, right). Samples were excited with a 10 ns 417 nm light pulse (0.6 mJ/pulse), and the transient absorption was monitored at 470 nm.

value measured for **1** in fluid water ( $k_q = 1.8 \times 10^9 \text{ M}^{-1} \text{ s}^{-1}$ ). The experiments for **1** in fluid water solution produced a  $K_D$  value of  $89\,600 \text{ M}^{-1}$ . Thus, the time scale for  $O_2$  diffusion may be slower in the water–PVG heterogeneous system compared to fluid solution. By analogy, Thomas et al. used the Einstein equation  $\langle x^2 \rangle = 6Dt$  to suggest that nitromethane quenching of excited anthracene is  $10^3$ – $10^6$  times slower in a zeolite compared to fluid solution (cf.  $10^{-8}$ – $10^{-11} \text{ cm}^2 \text{ s}^{-1}$ ).<sup>30</sup> In our system, the smaller  $K_D$  values in the heterogeneous samples may result from reduced directions for access of  $O_2$  to  $^3\mathbf{1}^*\text{-ads}$  compared to  $O_2$  to  $^3\mathbf{1}^*$  in fluid solution.

The amplitude of the time-resolved absorbance at 470 nm was examined (Figure S1, Supporting Information). Because a

ground-state adduct equilibrium of **1-ads** and O<sub>2</sub> would be expected to have a different absorption spectrum from **1-ads**, an absorption decrease can point to a static quenching mechanism. Figure 6 (right-hand Y-axis) shows the plot of  $\alpha_0/\alpha$  vs [O<sub>2</sub>] was linear over the O<sub>2</sub> concentration examined and followed the equation  $\alpha_0/\alpha = 1 + K_S[\text{O}_2]$ . For **1-ads**, the adduct formation constant  $K_S$  ranged from 510 to 1310 M<sup>-1</sup>. The amplitude of the absorption at 470 nm decreased slightly (by about 0.1) with increasing O<sub>2</sub> concentrations. Thus, we estimate that ~10% of the quenching occurred via the static quenching mechanism (colored blue in Scheme 2). The above Stern–Volmer analysis suggested that quenching of **31\***-**ads** by O<sub>2</sub> is primarily dynamic, the route colored black in Scheme 2. In the absence of aggregation, photosensitizers in homogeneous fluid solution are often quenched dynamically by O<sub>2</sub>.<sup>21</sup>

#### 4. Conclusion

In isotropic heterogeneous media static and dynamic quenching modes can operate. In the present case, we have evidence that O<sub>2</sub> quenches triplet **1\*** at the water–PVG interface primarily by a dynamic quenching mechanism. The contribution from static quenching remains low even with loadings of **1** onto PVG that varied by 100-fold. The O<sub>2</sub> quenching constants  $k_q$  for **31\***-**ads** were only 3–4 times smaller than for **31\*** itself, suggesting the heterogeneous system is capable of generating singlet oxygen for its use as a reagent in the surrounding aqueous solution.

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**Supporting Information Available:** Time-resolved absorbance of **31\***-**ads** in the presence of increasing concentrations of oxygen. This material is available free of charge via the Internet at <http://pubs.acs.org>.

#### References and Notes

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